

Mata Kuliah : Dinamika Struktur & Pengantar Rekayasa Kegempaan

Kode : TSP - 302

SKS : 3 SKS

Multi Degree of Freedom System Free Vibration

Pertemuan – 6, 7



• TIU:

- Mahasiswa dapat menjelaskan tentang teori dinamika struktur.
- Mahasiswa dapat membuat model matematik dari masalah teknis yang ada serta mencari solusinya.

• TIK :

Mahasiswa mampu mendefinisikan derajat kebebasan, membangun persamaan gerak sistem MDOF



Sub Pokok Bahasan :

- Penentuan derajat kebebasan
- Properti matrik kekakuan, massa dan redaman



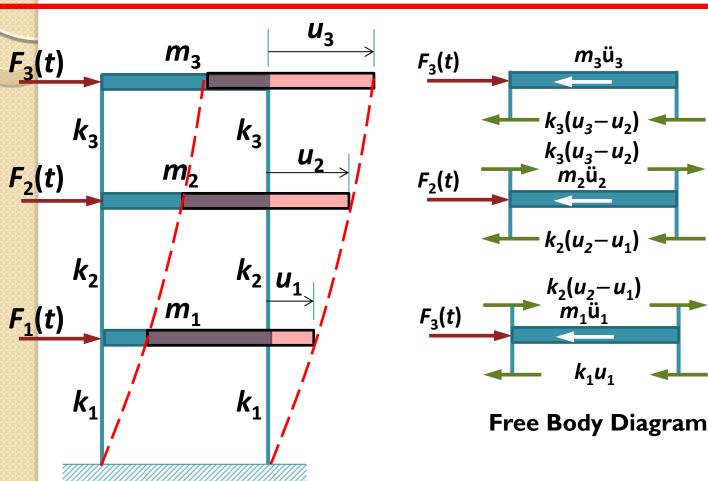
- Structures cannot always be described by a SDoF model
- In fact, structures are continuous systems and as such possess an infinite number of DoF
- There are analytical methods to describe the dynamic behavior of continuous structures, which are rather complex and require considerable mathematical analysis.
- For practical purposes to study Multi Degree of Freedom System (MDoF), we shall consider the multistory shear bulding model.



- A shear building may be defined as a structure in which there is no rotation of a horizontal section at the level of the floors
- Assumption for shear building
 - The total mass of the structure is concentrated at the levels of the floors
 - The slabs/girders on the floors are infinitely rigid as compared to the columns
 - The deformation of the structure is independent of the axial forces present in the columns



Shear Building w/ 3 DoF





By equating to zero the sum of the forces acting on each mass :

$$m_{1}\ddot{u}_{1} + k_{1}u_{1} - k_{2}(u_{2} - u_{1}) - F_{1}(t) = 0$$

$$m_{2}\ddot{u}_{2} + k_{2}(u_{2} - u_{1}) - k_{3}(u_{3} - u_{2}) - F_{2}(t) = 0$$

$$m_{3}\ddot{u}_{3} + k_{3}(u_{3} - u_{2}) - F_{3}(t) = 0$$

$$(1)$$

Eq. (I) may conveniently be written in matrix notation as:

$$[M]\{\ddot{u}\} + [K]\{u\} = \{F\}$$

Where :

And :

$$\{u\} = \begin{cases} u_1 \\ u_2 \\ u_3 \end{cases} \qquad \{\ddot{u}\} = \begin{cases} \ddot{u}_1 \\ \ddot{u}_2 \\ \ddot{u}_3 \end{cases} \qquad \{F\} = \begin{cases} F_1(t) \\ F_2(t) \\ F_3(t) \end{cases}$$

Natural Frequencies and Normal Modes

EoM for MDoF system in free vibration is :

$$[M]{\ddot{u}} + [K]{u} = \{0\}$$

Solution of Eq. (3) is :

$$\{u\} = \{\phi_n\}q_n(t) = \{\phi_n\}(A_n\cos\omega_n t + B_n\sin\omega_n t)$$
 (4)

The substitution of Eq. (4) into Eq. (3) gives :

$$[K]\{\phi_n\} - \omega_n^2 [M]\{\phi_n\}] q_n(t) = \{0\}$$
(5)



- The formulation of Eq (5) is known as an eigenvalue problem.
- To indicate the formal solution to Eq. (5), it is rewritten as :

$$[K] - \omega_n^2 [M] \{\phi_n\} = \{0\}$$
 (6)

Eq (6) has non trivial solution if:

$$\left[\left[\mathbf{K} \right] - \omega_n^2 \left[M \right] = 0$$
 (7)

• Eq.(7) gives a polynomial equation of degree n in ω_n^2 , known as characteristic equation of the system

A vibrating system with n DoF has n natural vibration frequencies ω_n , arranged in sequence from smallest to largest



- The n roots of Eq. (7) determine the n natural frequencies ω_n of vibration.
- These roots of the characteristic equation are also known as eigenvalues
- When a natural frequency ω_n is known, Eq.(6) can be solved for the corresponding vector ϕ_n
- There are n independent vectors ϕ , which are known as natural modes of vibration, or natural mode shape of vibration
- These vectors are also known as eigenvector.

Example I

Determine :

The natural frequencies and corresponding modal shapes

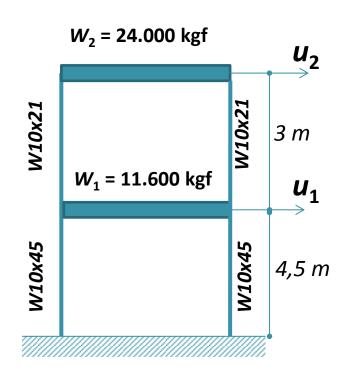
$$W10x21 \rightarrow I = 4,900 \text{ cm}^4$$

$$W10x45 \rightarrow I = 10,300 \text{ cm}^4$$

$$E = 200.000 MPa$$

$$\omega_1 = 10,76$$
 $\omega_2 = 38,28$

$$\begin{bmatrix} \phi \end{bmatrix} = \begin{bmatrix} \phi_{11} & \phi_{12} \\ \phi_{21} & \phi_{22} \end{bmatrix} = \begin{bmatrix} 0,6809 & -3,0385 \\ 1 & 1 \end{bmatrix}$$



Modal & Spectral Matrices

 The N eigenvalues, N natural frequencies, and n natural modes can be assembled compactly into matrices

$$\left[\Omega^2\right] = \begin{bmatrix} \omega_1^2 & & & \\ & \omega_2^2 & & \\ & & \ddots & \\ & & \omega_n^2 \end{bmatrix}$$



Spectral matrix

Orthogonality of Modes

The natural modes corresponding to different natural frequencies can be shown to satisfy the following orthogonality conditions:

$$\{\phi_n\}^T [K] \{\phi_r\} = 0$$

$$\{\phi_n\}^T [M] \{\phi_r\} = 0$$
 When $\omega_n \neq \omega_r$



Free Vibration Response: Undamped System

EoM for MDoF Undamped Free Vibration is :

$$[M]\{ii\} + [K]\{u\} = \{0\}$$
(8)

- The differential equation (8) to be solved had led to the matrix eigenvalue problem of Eq.(6).
- The general solution of Eq. (8) is given by a superposition of the response in individual modes given by Eq. (4).

$$\{u\} = \{\phi_1\}q_1 + \{\phi_2\}q_2 + \{\phi_3\}q_3 + \dots + \{\phi_n\}q_n\}$$
 (9)

• Or
$$\{u\} = \sum_{n=1}^{N} \phi_n q_n$$
 $\{u\} = [\phi]\{q\}$ (10)



• Premultiply Eq. (10) with $\{\phi_n\}^T[M]$:

$$\{\phi_n\}^T [M] \{u\} = \{\phi_n\}^T [M] [\phi] \{q_n\}$$

$$= \{\phi_n\}^T [M] \{\phi_1\} q_1 + \{\phi_n\}^T [M] \{\phi_2\} q_2 + \dots$$

$$+ \{\phi_n\}^T [M] \{\phi_N\} q_N$$

$$(11)$$

Regarding the orthogonality conditions:

$$\{\phi_n\}^T [M] \{u\} = \{\phi_n\}^T [M] \{\phi_n\} q_n$$

• From which :

$$q_{n} = \frac{\{\phi_{n}\}^{T}[M]\{u\}}{\{\phi_{n}\}^{T}[M]\{\phi_{n}\}} \quad \text{also} \quad \dot{q}_{n} = \frac{\{\phi_{n}\}^{T}[M]\{\dot{u}\}}{\{\phi_{n}\}^{T}[M]\{\phi_{n}\}} \quad (11)$$



If Eq. (8) is premultiplied by the transpose of the nth modeshape vector $\{\phi_n\}^T$, and using Eq.(10) with its second derivative it becomes:

$$\{\phi_n\}^T [M] \{\phi_n\} \ddot{q}_n + \{\phi_n\}^T [K] \{\phi_n\} q_n = 0$$

$$\mathbf{K}_n$$

$$\mathbf{K}_n$$

$$(11)$$

Generalized mass Generalized stiffness

Eq.(11) is an EoM from SDoF Undamped System for mode n, which has solution:

$$q_n(t) = q_n(0)\cos\omega_n t + \frac{\dot{q}_n(0)}{\omega_n}\sin\omega_n t$$
 (12)



Solution for Eq. (8) becomes :

$$\{u\} = \sum_{n=1}^{N} \phi_n \left[q_n(0) \cos \omega_n t + \frac{\dot{q}_n(0)}{\omega_n} \sin \omega_n t \right]$$
 (13)

- The procedure described above can be used to obtain an independent SDoF equation for each mode of vibration of the undamped structure.
- This procedure is called the mode-superposition method



Example 2

 Based on data from Example I, and the following initial condition, for each DoF plot the time history of displacement, regarding the Ist, and 2nd mode contribution.

$$u_{(t=0)} = \begin{cases} 6 \cdot 10^{-3} \\ 8 \cdot 10^{-3} \\ 10 \cdot 10^{-3} \end{cases} m \qquad \dot{u}_{(t=0)} = \begin{cases} 0 \\ 0,20 \\ 0 \end{cases} m/s$$

Free Vibration Response: Damped System

EoM for MDoF Damped Free Vibration is :

$$[M]\{\ddot{u}\} + [C]\{\dot{u}\} + [K]\{u\} = \{0\}$$
 (14)

Using Eq. (10) and its derivative :

$$[M][\phi]{\ddot{q}} + [C][\phi]{\dot{q}} + [K][\phi]{q} = \{0\}$$
 (15)

• Premultiply Eq. (15) with $\{\phi\}^T$:



Eq.(16) is an EoM from SDoF Damped System for mode n, which has solution:

$$q_n(t) = e^{-\xi_n \omega_n t} \left[q_n(0) \cos \omega_{nD} t + \frac{\dot{q}_n(0) + \xi_n \omega_n q_n(0)}{\omega_{nD}} \sin \omega_{nD} t \right]$$

$$(17)$$

 The displacement response of the system is then obtained by substituting Eq. (17) in Eq. (10)

$$u(t) = \sum_{n=1}^{N} \phi_n e^{-\xi_n \omega_n t} \left[q_n(0) \cos \omega_{nD} t + \frac{\dot{q}_n(0) + \xi_n \omega_n q_n(0)}{\omega_{nD}} \sin \omega_{nD} t \right]$$
(18)